

Critical Point Scans at RHIC Full Energy

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“The QCD Critical Point Workshop”

INT, Aug 13, 2008

n_B/s is a “figure of merit” for any isentropic trajectory

How much experimental control do we have?

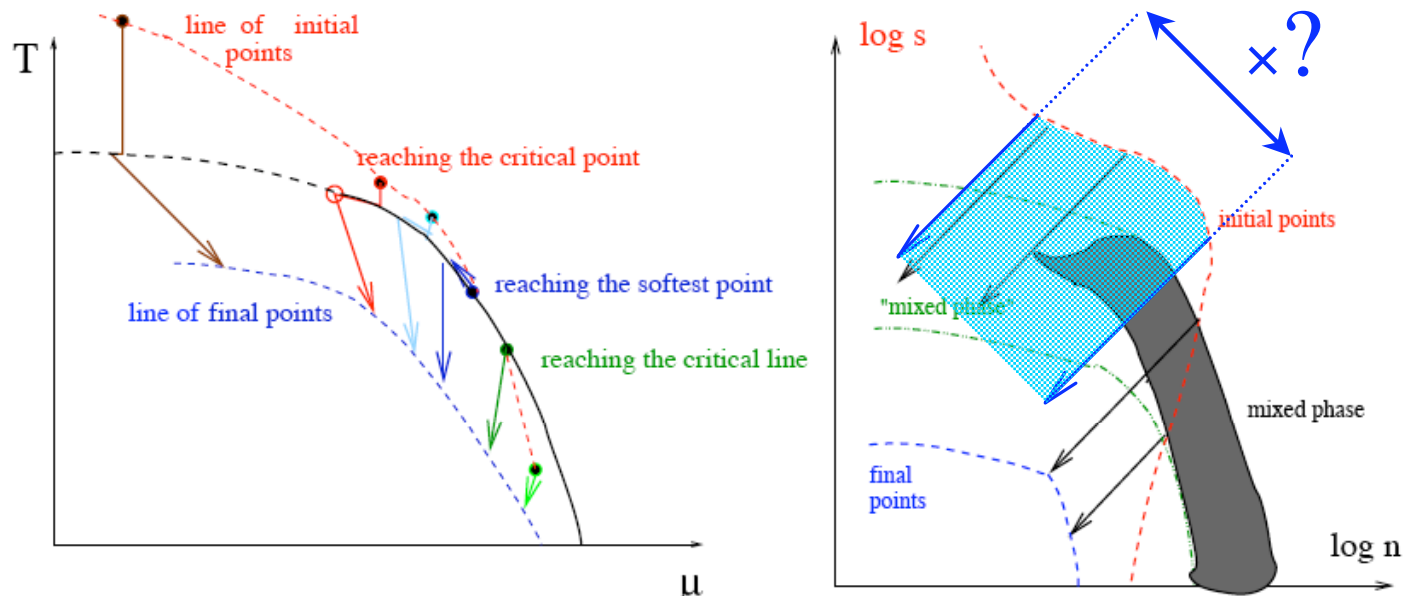
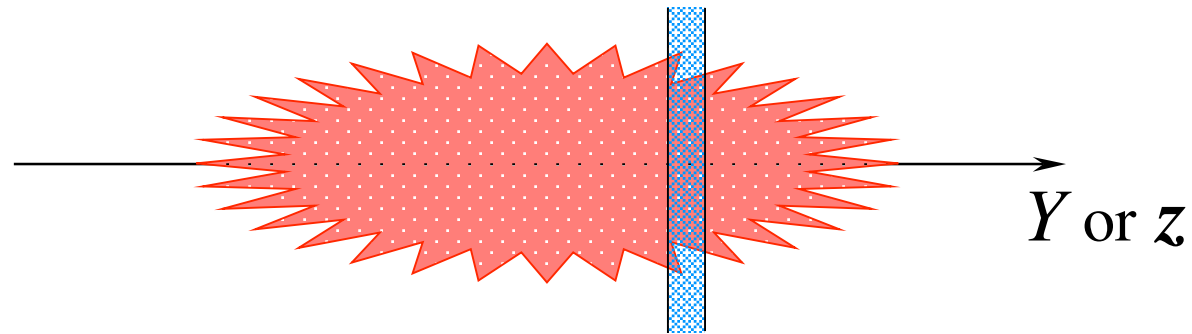


Figure 2: Schematic view of the cooling paths on (a) $T - \mu$ and (b) $\log(s) - \log(n)$ diagrams. In (a) those paths are zigzag-shaped lines with arrows, extended from the (red dashed) line of the initial points, to (blue dashed) line of the kinetic freezout. In (b) the same lines are a set of parallel straight lines.

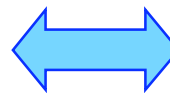
“All thermodynamics is local”



Scaling Solution:

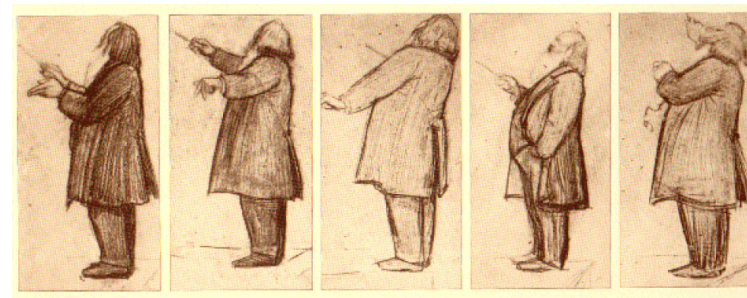
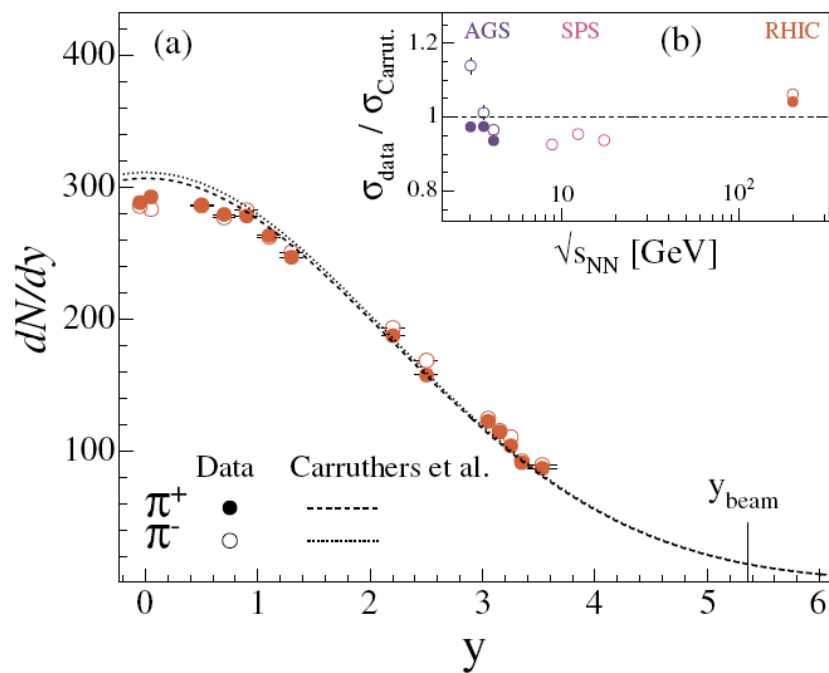
(not restricted to Bjorken hydrodynamics)

Fluid at/near
position z at
some time



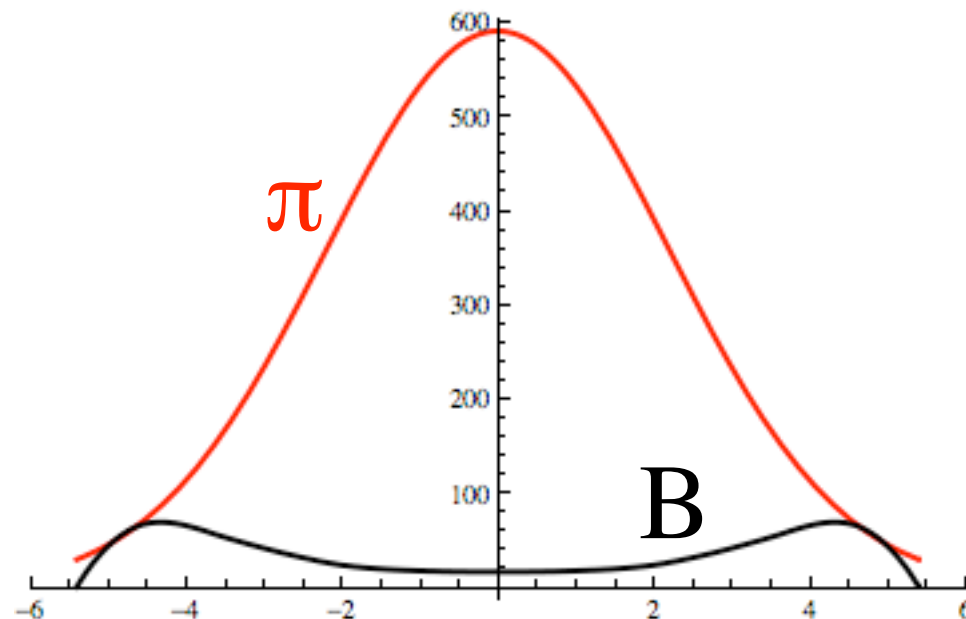
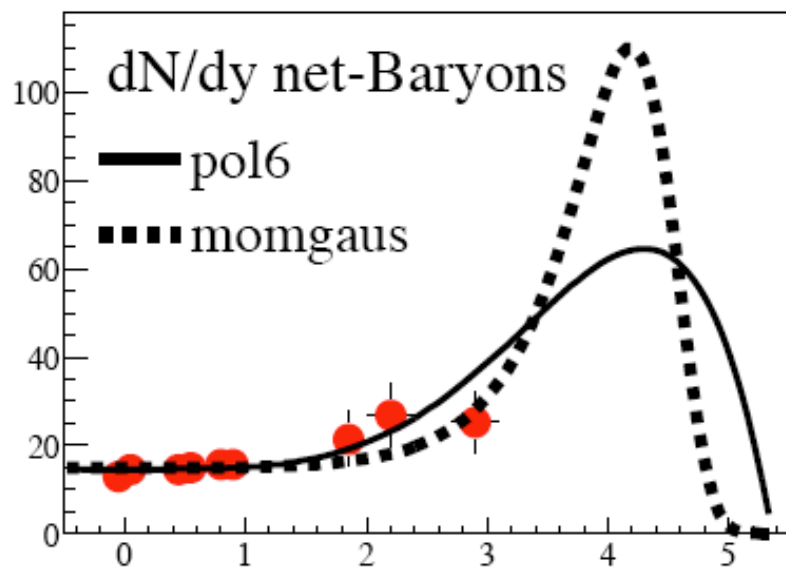
Fluid velocity at/near
some rapidity Y at
that time

$$\left(\frac{n_B}{s}\right)_{\text{Initial}} = \left(\frac{n_B}{s}\right)_{\text{Final}} \approx C \frac{n_B}{n_\pi} = C \frac{dN^B / dY}{dN^\pi / dY} \equiv C \left(\frac{B}{\pi}\right)$$



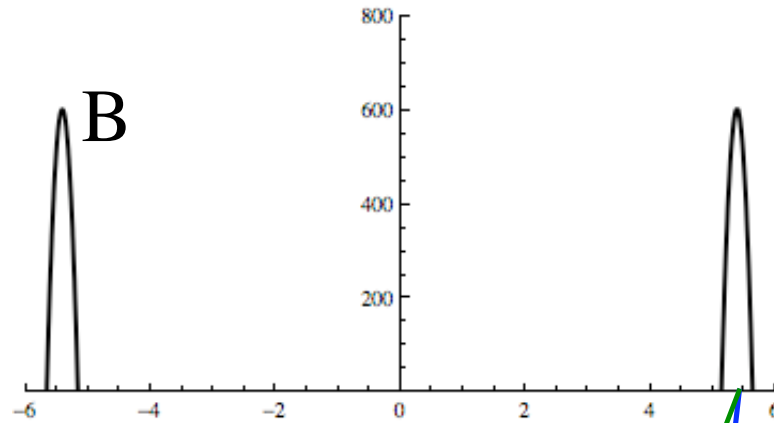
BRAHMS PRL 94 (2005) 162301

BRAHMS PRL 93 (2004) 102301

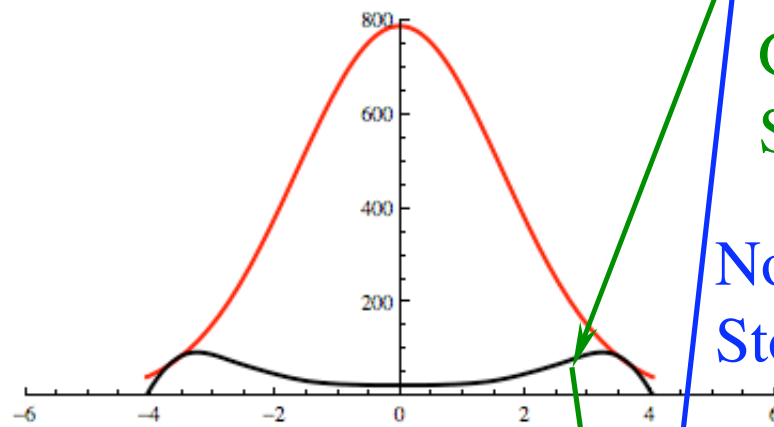


RHIC Au+Au 200 GeV Central

Initial



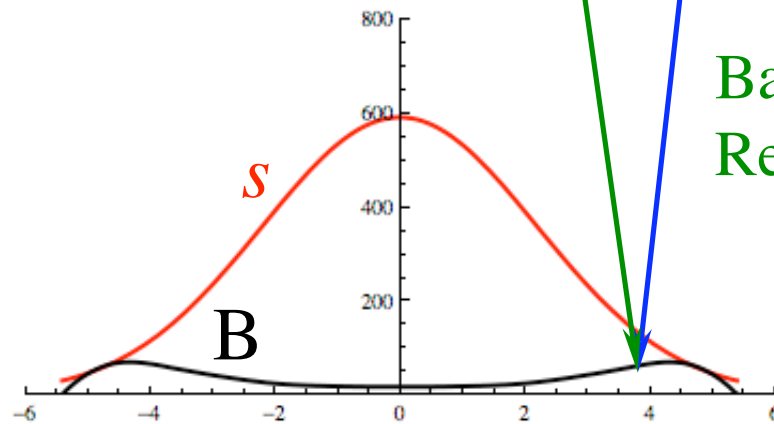
Early Hydro



Greater Baryon Stopping

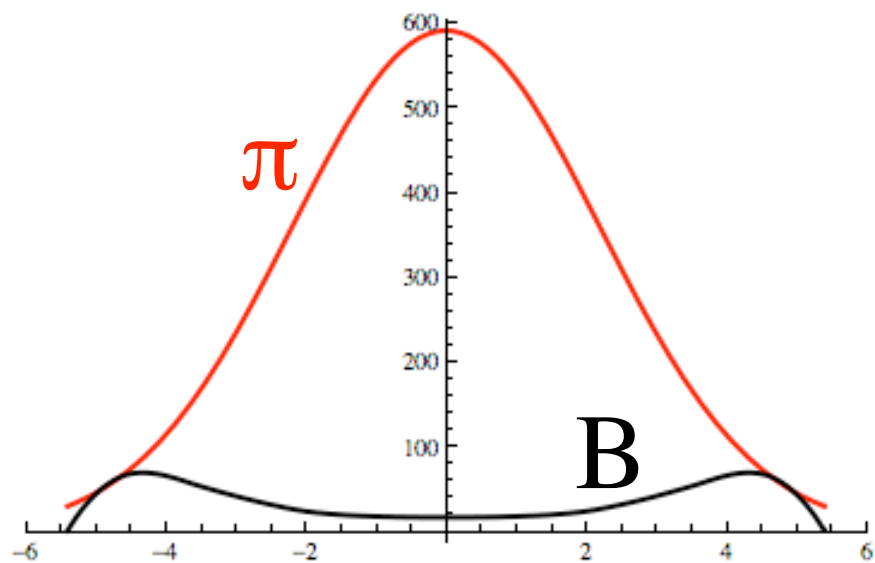
Nominal Baryon Stopping

Late Hydro



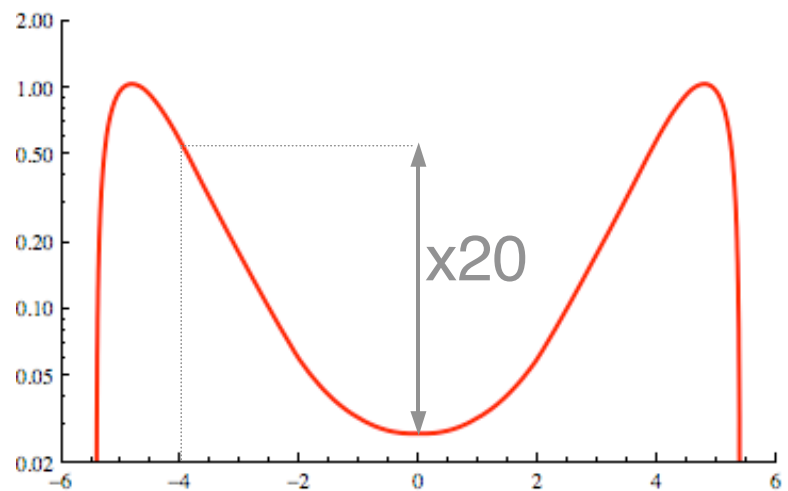
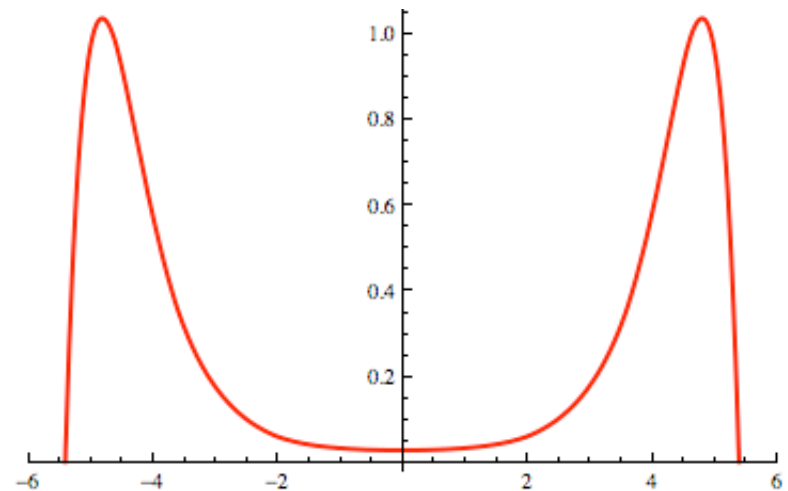
Baryon Re-Acceleration!

Rapidity Scan

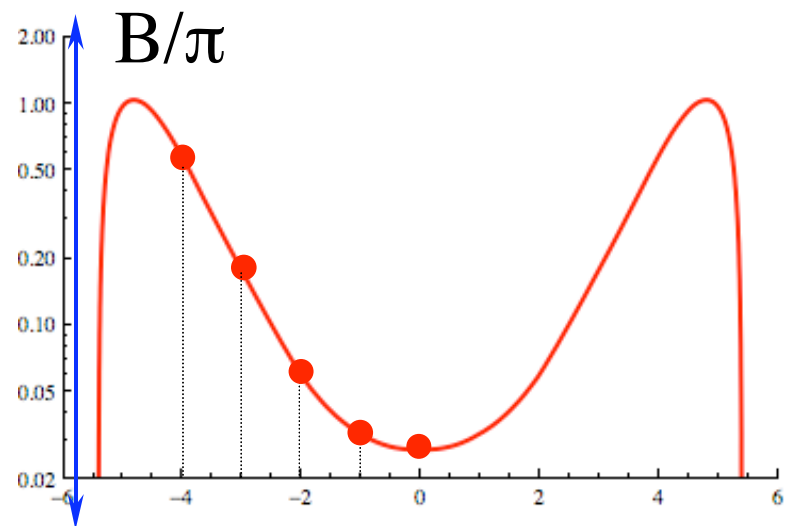
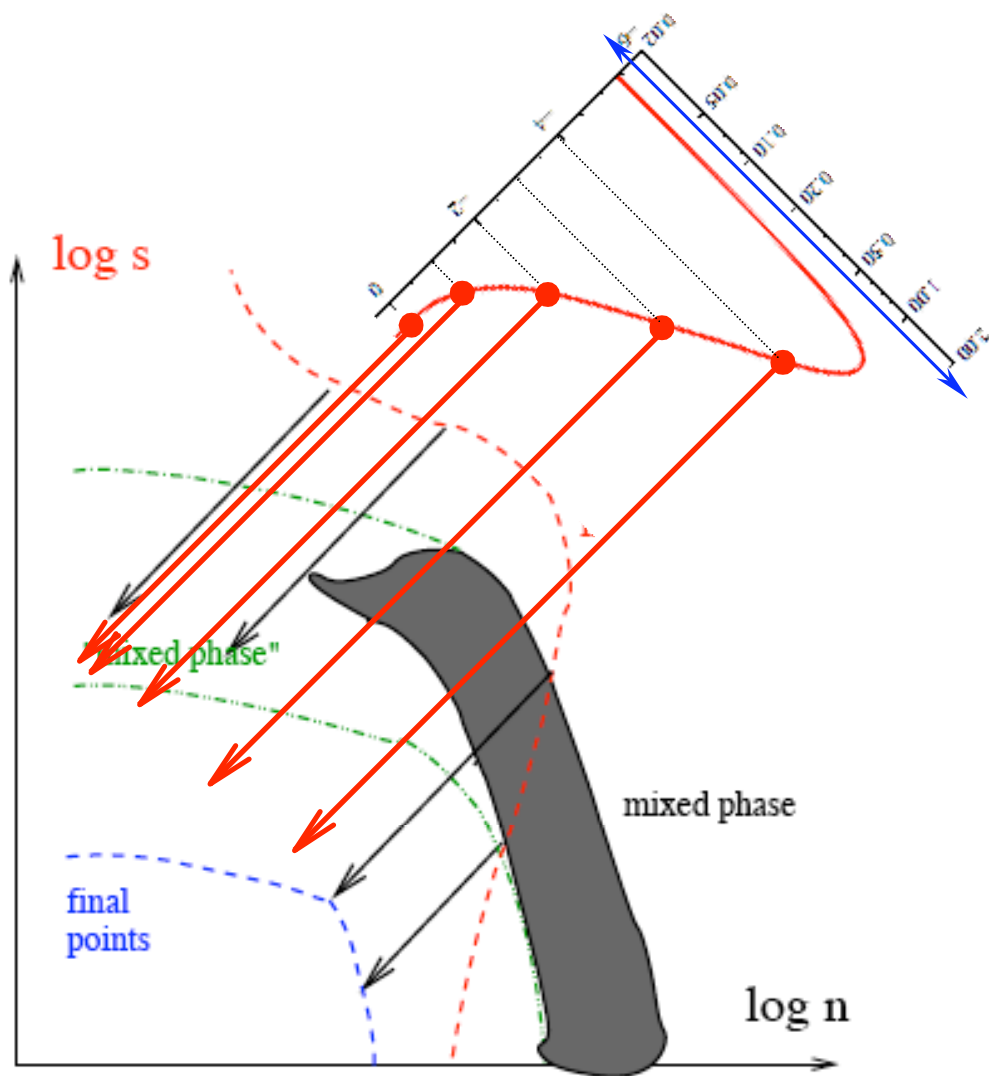


RHIC Au+Au 200 GeV Central

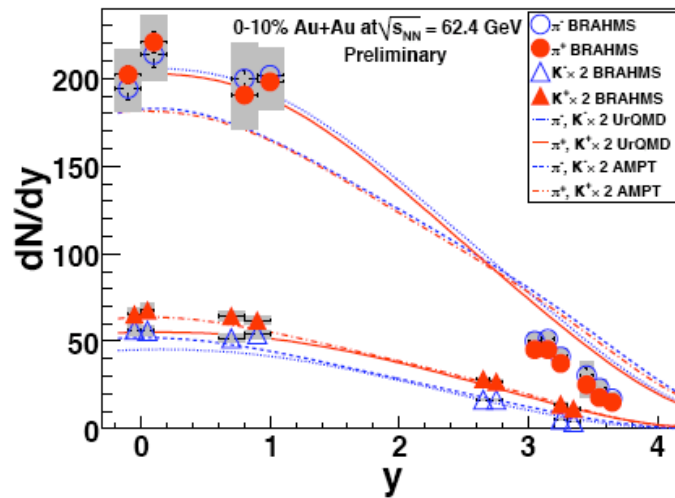
B/π



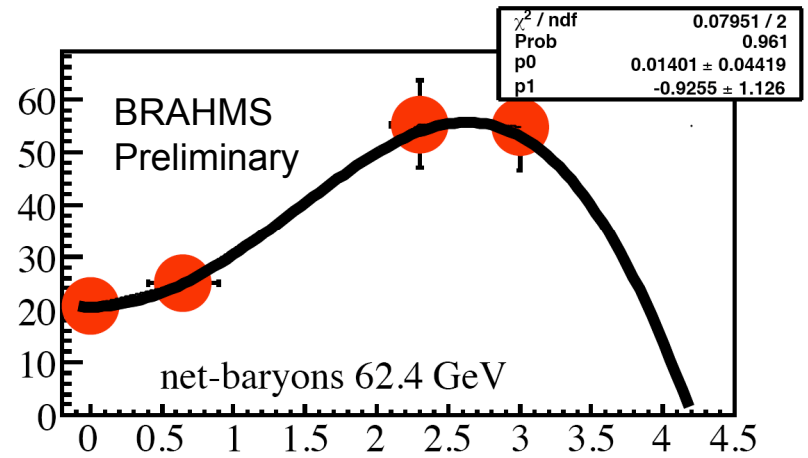
Every collision contains a scan!



62

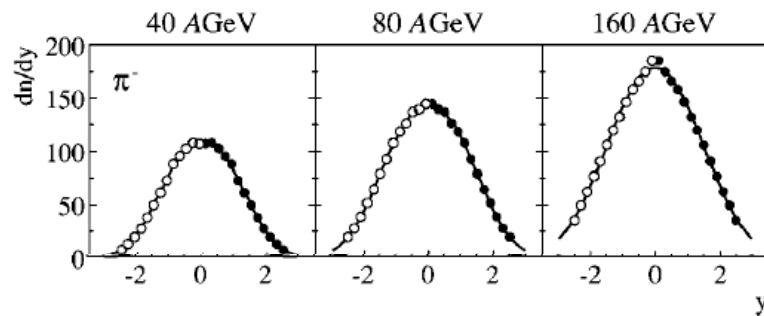


BRAHMS I.C. Arsene QM2008

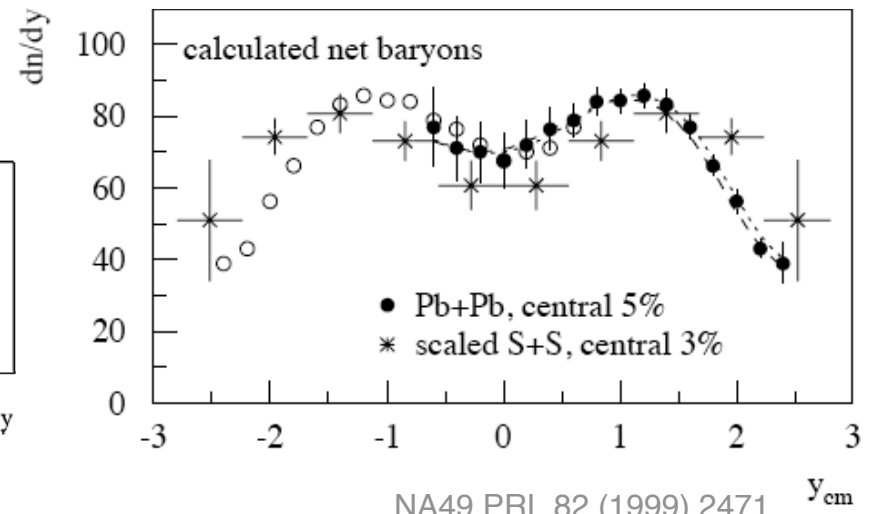


BRAHMS H.H. Dalsgaard QM2008

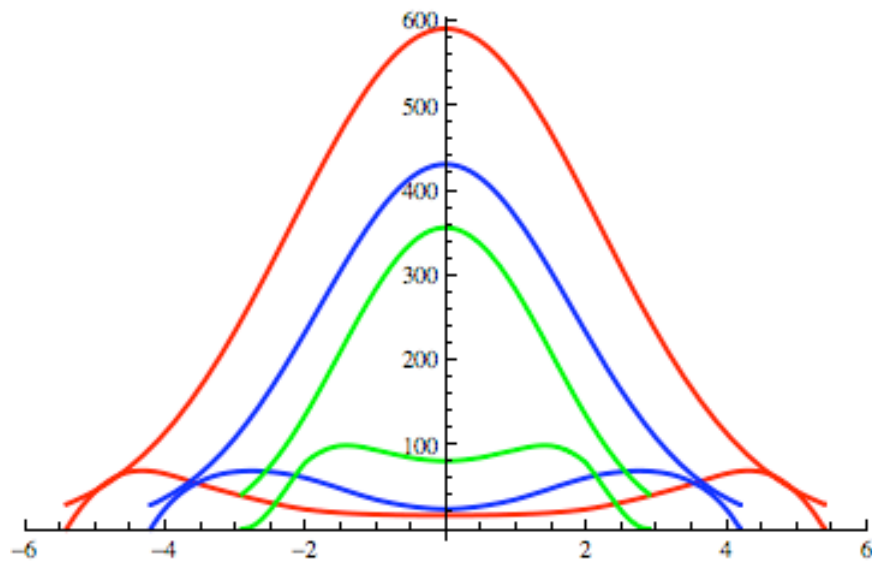
17



NA49 PRC 56 (2002) 054902

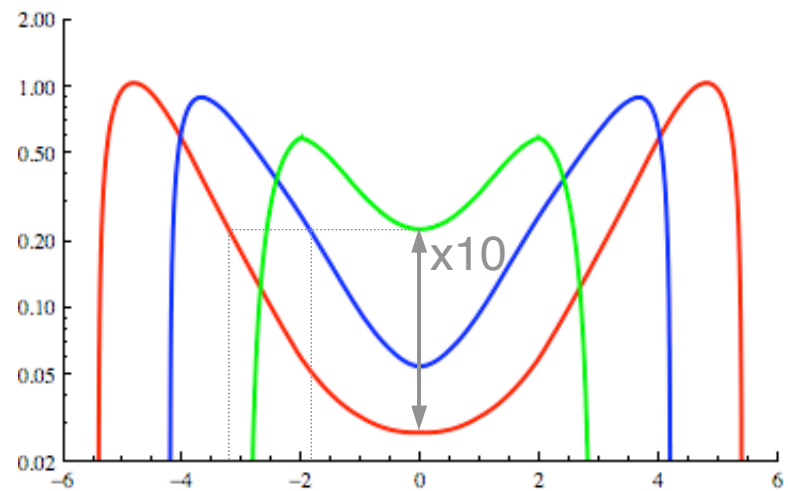
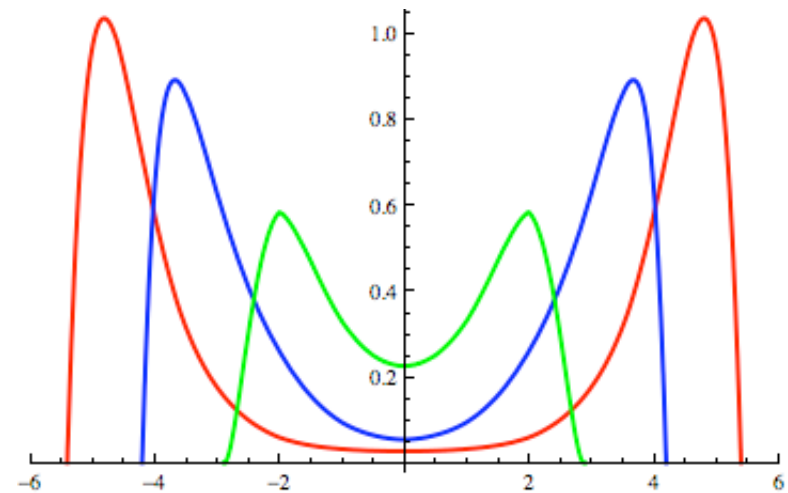
NA49 PRL 82 (1999) 2471 y_{cm}

Beam Energy Scan



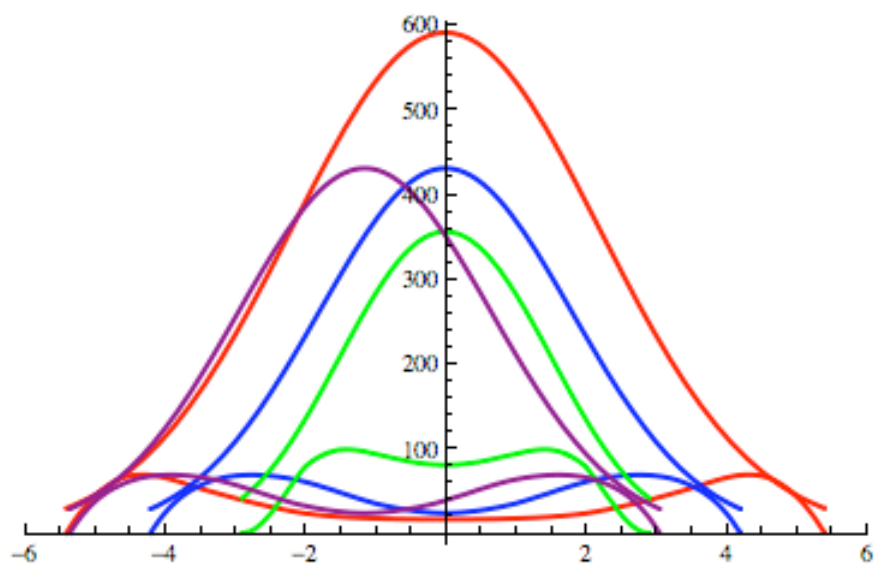
$$\sqrt{s_{NN}} = \text{200} \text{ 62} \text{ 17}$$

B/π



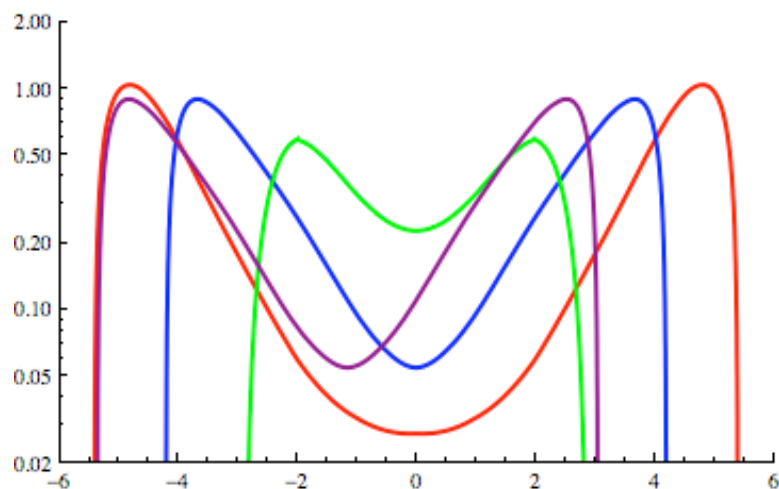
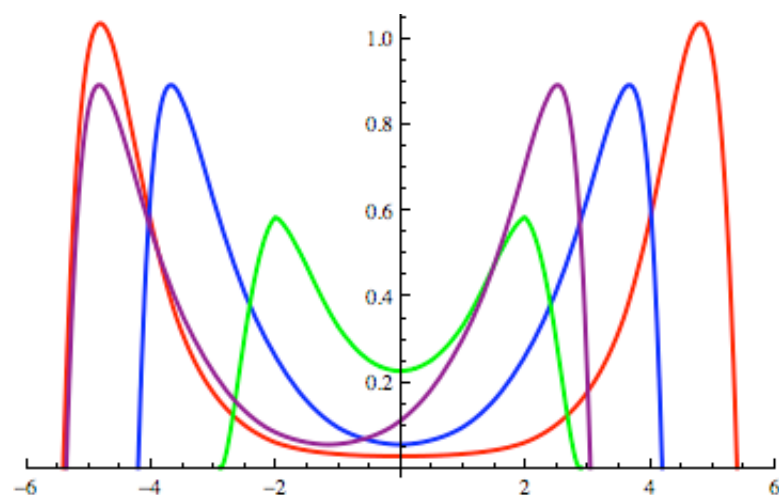
Asymmetric Beam Energies

A 100 AGeV + 10 AGeV collision
has $\sqrt{s_{\text{NN}}} \approx 60$ GeV and $\Delta Y_{\text{CMS}} \approx 1.15$

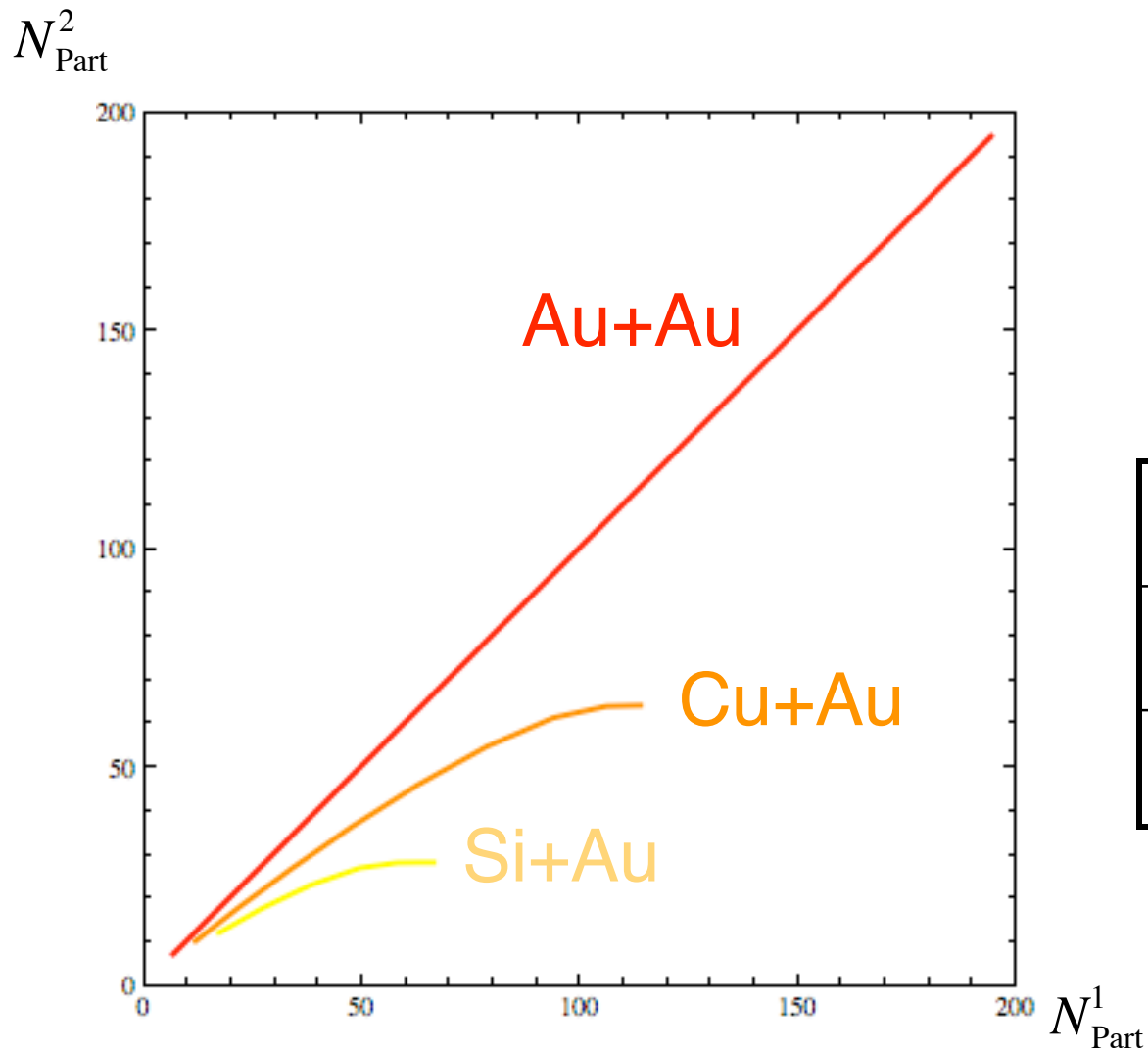


200 100+10 62 17

B/π



Asymmetric Ion Collisions: N_{Part}

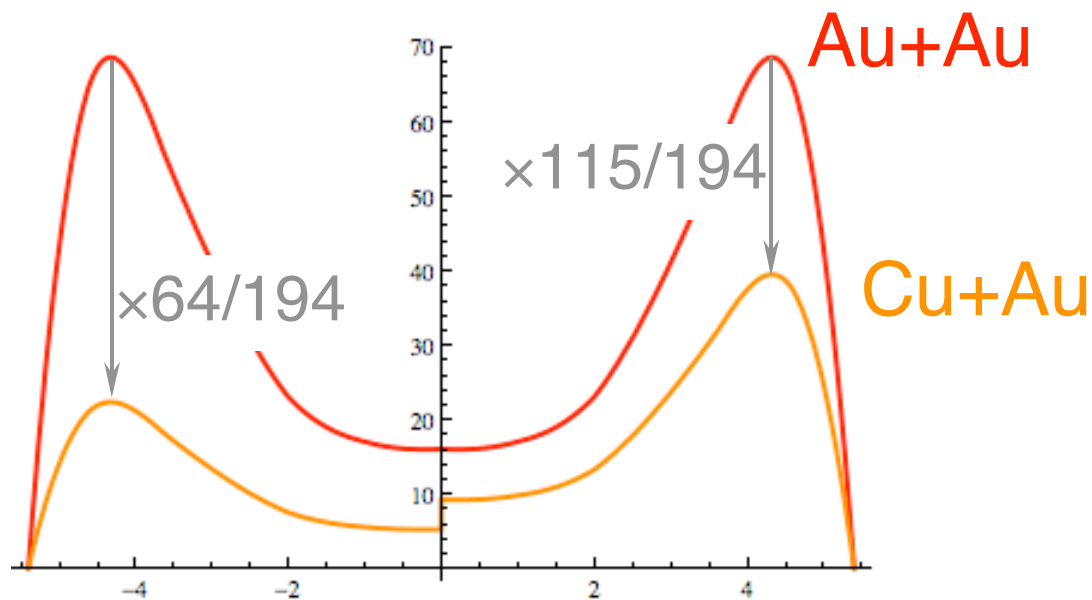


Central $b=0$

	N_{Part}^1	N_{Part}^2
Au+Au	194	194
Cu+Au	64	115
Si+Au	28	67

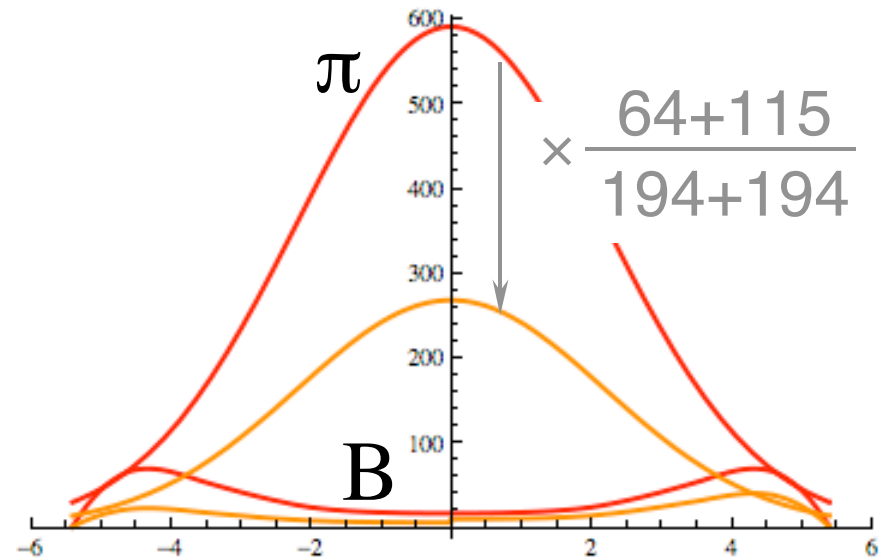
Asymmetric Ion Collisions: dN/dY

$$dN_B/dY$$



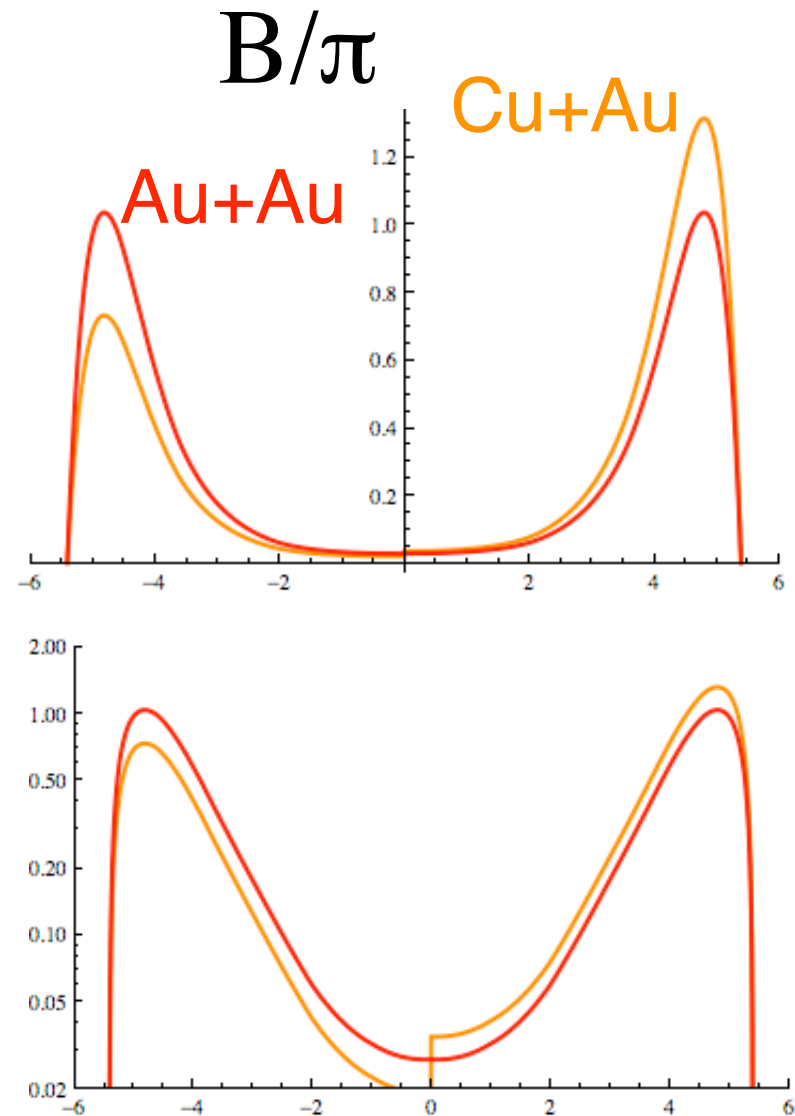
	N_{Part}^1	N_{Part}^2
Au+Au	194	194
Cu+Au	64	115

Assume: Each dN_B/dY scales as N_{Part} ,
and dN_π/dY scales as $N_{\text{Part}}^1 + N_{\text{Part}}^2$

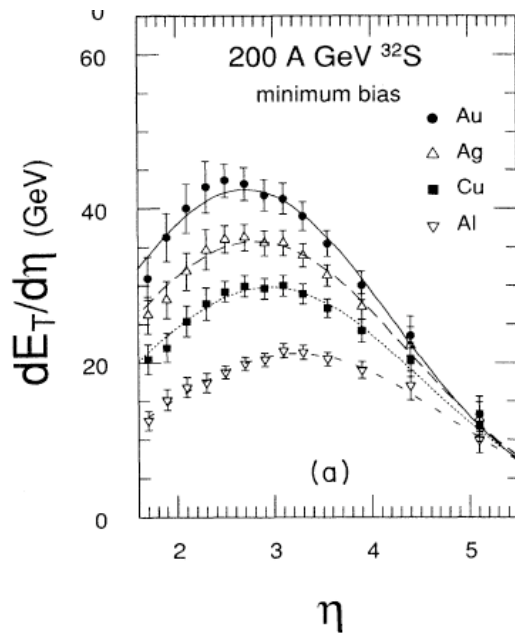


Asymmetric Ion Collisions: B/π

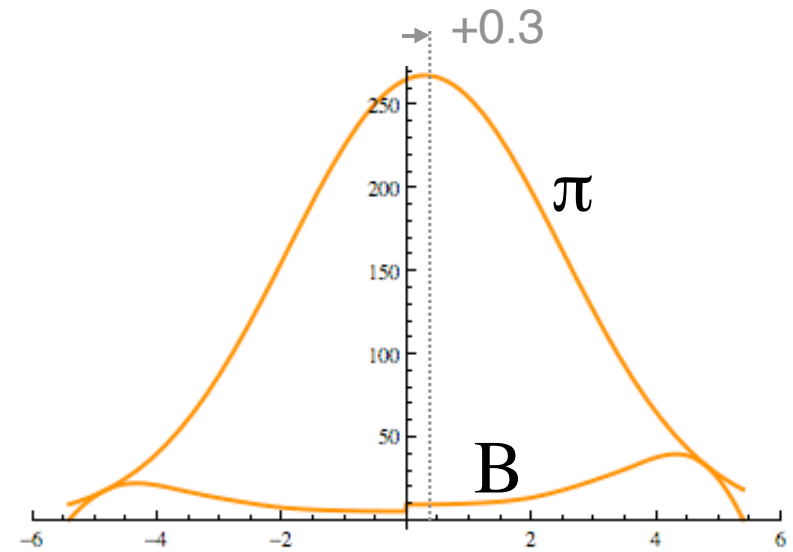
Under simple N_{Part} scaling, if $f \equiv N_{\text{Part1}} / N_{\text{Part2}}$ then B/π changes by a factor of $2/(1+f)$ relative to symmetric collisions.



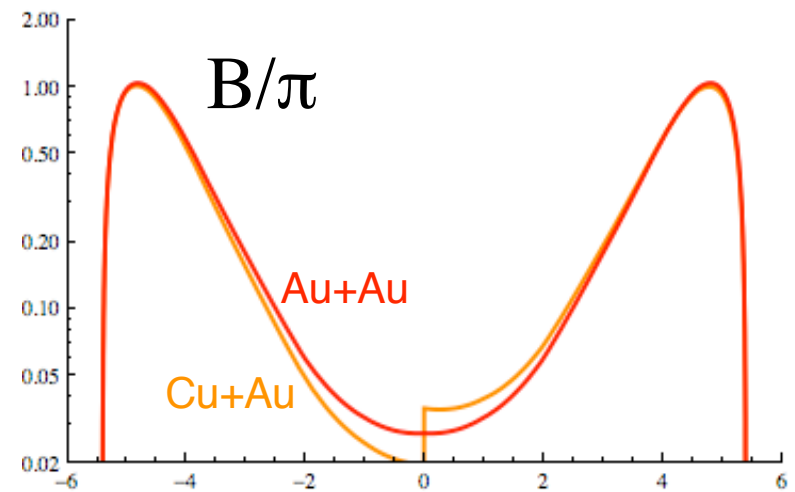
Shifting π 's relative to B's



WA80 PRC 44
(1991) 2736



	N_{Part}^1	N_{Part}^2	ΔY_{Part}
Au+Au	194	194	0
Cu+Au	64	115	0.3
Si+Au	28	67	0.44



Summary:

1. For isentropic evolution, n_B/s is conserved for a fluid element. We can approximate n_B/s for a fluid element by B/π at some rapidity Y in the final state.
2. At full RHIC energy B/π is a strong function of rapidity, varying over a factor of $\times 20$ from $Y=0$ to $Y=4$. Effectively every collision contains a scan in n_B/s .
3. Beam energy scans from top RHIC to top SPS vary B/π by about a factor of $\times 10$ at mid-rapidity. Asymmetric beam energy collisions can shift the B/π pattern significantly.
4. Asymmetric ion collisions at RHIC are of fundamental interest for many reasons but don't provide much leverage in B/π , under simple assumptions.